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# High-sensitivity refractive index sensor based on Ge–Sb–Se chalcogenide microring resonator

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### ABSTRACT

In this article, we report a liquid sensing device based on a Ge–Sb–Se microring resonator (MRR) operating at a wavelength of 1550 nm. The transmission loss of the waveguide is 4.3 dB/cm, and the intrinsic Q factor of the MRR is 7.74  $\times$  10<sup>4</sup>, with an extinction ratio of approximately 40 dB. By detecting the different concentrations of NaCl solutions, we observe that the sensitivity of our proposed device is approximately 123 nm/RIU, and the intrinsic limit of detection is approximately 3.24  $\times$  10<sup>-4</sup> RIU. Compared with the other materials, such as silicon, Si<sub>3</sub>N<sub>4</sub>, or polymer, the moderate refractive index (RI) of Ge<sub>28</sub>Sb<sub>12</sub>Se<sub>60</sub> glass allows it to obtain a high sensitivity and a high Q factor, while maintaining a strong confinement on the light field to avoid suffering from bending loss under a small bending radius, with a compact structure.

#### 1. Introduction

As an important part of information technology, sensing technology shoulders the important task of obtaining external information. Together with communication and computer technologies, they form the "senses", "nerve" and "brain" of information system, to jointly support the rapid development of the information society [1,2]. Among different types of sensors, optical sensor based on whispering gallery mode (WGM) provides low limit of detection (LOD) and high sensitivity, due to its high quality factor and small mode field volume [3–6]. To date, several types of WGM sensors with different structures have been reported, such as microrings, microdisks, microspheres, and microtoroids [7–9]. Among which, the WGM sensors based on microring resonator (MRR) have attracted much interest due to their high sensitivity to the change of external refractive index (RI) and the ability to work under single-mode condition [10–12]. As a optical resonant microcavity with compact structure, the light field can be confined to the order of micronano scale without using additional reflective mirrors. The cyclic resonance of the light field in such a small area can significantly enhance the interaction with external analyte, which will greatly improve the sensitivity and *LOD* of the sensor. Meanwhile, the fabrication of the microring is fully compatible with the CMOS semiconductor micro-nano manufacturing process, which is suitable to build a fully functional "lab on chip" system.

Many materials, such as lithium niobite [13], silicon nitride [14], III–V group semiconductors [15], polymers [16], and chalcogenide glasses (ChGs) [17–19], have been used to fabricate various photonic devices for sensing applications. ChGs are amorphous materials formed by covalent bonding of elements such as sulfur, selenium, and tellurium in Group VIA, with a certain amount of other metal or non-metal elements [20,21]. Compared with oxide counterparts, chalcogenides have higher glass density and contain atoms with high polarizability, which leads to higher linear and nonlinear RIs. In addition, because the atomic mass of chalcogenide elements (S, Se, Te) is heavier than that of oxygen, they have lower phonon energy (200–400 cm<sup>-1</sup>), which makes the transparency window can be extended up to 12, 16 and 25  $\mu$ m, for sulfides, selenides and tellurides respectively [22,23]. Such a wide infrared transparency window covers the characteristic absorption

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Received 11 March 2021; Received in revised form 19 May 2021; Accepted 21 May 2021 Available online 28 May 2021 1350-4495/© 2021 Elsevier B.V. All rights reserved. bands for many biochemical molecules, which makes ChGs a very promising platform for infrared sensing applications.

Meanwhile, the RI of ChGs is higher than that of materials such as silicon nitride or polymers, which enables ultra-compact photonic devices without suffering from large bending losses. Furthermore, the amorphous property of ChGs enable us to efficiently deposit large-area and high-quality thin films on almost any substrate at a very low cost without considering lattice mismatch. Several ChG materials such as As<sub>2</sub>S<sub>3</sub> [24], As<sub>2</sub>Se<sub>3</sub> [25], and Ge<sub>11.5</sub>As<sub>24</sub>Se<sub>64.5</sub> glasses [26], have been used to make photonic devices for various sensing applications. However, the extreme harm of Arsenic to the natural environment and human beings will greatly limit their application. At this time, Arsenicfree and selenium-based ChGs provide an alternative option. Several studies have illustrated that the thermal and chemical stabilities of Ge-Sb-Se glasses are comparable to those of As-Se glasses, with a relatively large glass-formation region and excellent mid-infrared transparency [27,28]. Moreover, according to Miller's law, the third-order nonlinearity of selenium-based glass can be improved after replacing arsenic with antimony, because the stronger cationic polarizability of antimony [29,30]. Recently, these excellent properties of Ge-Sb-Se glasses have attracted interest in various fields, such as supercontinuum generation [31,32], wavelength conversion [33], all-optical signal processing, and sensing.

In this article, we report a compact  $Ge_{28}Sb_{12}Se_{60}$  MRR sensor fabricated on a standard Si/SiO<sub>2</sub> platform. By adopting a microring structure with the resonance enhancement effect, the refractive index sensitivity as high as 123 nm/RIU is obtained through sensing in aqueous solutions with different NaCl concentrations. We prove that chalcogenide materials have a great application prospect not only in the mid-infrared wavelength range but also in the near-infrared wavelength range.

#### 2. Microring resonator design

 $Ge_{28}Sb_{12}Se_{60}$  glass is selected as core materials, and silicon wafers with a 2 µm thick layer of silicon dioxide are used as a buffer to design MRR devices. The thermally grown oxide layer helps prevent light

leakage into the substrate due to the higher index contrast between the ChG and the silicon dioxide. Fig. 1(a) shows the schematic of the proposed sensor device, which consists of two bus waveguides and a ring resonator. Fig. 1(b) shows the cross-section view of the device, in which the waveguides and resonators have a height of 300 nm, and a vertical sidewall profile without losing generality is assumed. Fig. 1(c) shows the simulated neff of TE<sub>0</sub>, TE<sub>1</sub>, TE<sub>2</sub>, TE<sub>3</sub>, TM<sub>0</sub>, TM<sub>1</sub>, and TM<sub>2</sub> modes with different waveguide widths based on Finite-Difference Eigenmode (FDE) solver (MODE Solutions). The figure shows that the fundamental and first-order modes of TE/TM can be effectively supported when the width of the waveguide is greater than 1000 nm, but only the fundamental mode of TE/TM can be supported when the width of the waveguide is reduced below 750 nm. To avoid multiple modes, which may interfere with one another during propagation and will cause unwanted resonance peaks in the output spectrum, the width of the waveguide should be controlled below 750 nm. However, a very small size of the resonator will cause not only a high scattering loss due to the strong interaction between the light field and sidewall of the waveguide but also a challenge in device fabrication and characterization. Finally, the microring and the coupling waveguides are determined to have cross-sectional dimensions of 600 nm  $\times$  300 nm as a trade-off between the singlemode working condition, scattering loss, fabrication difficulty, and device robustness.

The fundamental mode profile for TE mode of the waveguide is shown in Fig. 2(a), in which the  $n_{eff}$  is calculated as approximately 2.002. By integrating the Poynting vector, about 80.386% of the power is confined in the waveguide core, and the effective area ( $A_{eff}$ ) of the mode field is calculated to be 0.3  $\mu$ m<sup>2</sup> using the following expression:

$$A_{eff} = \frac{\left(\iint |E(x,y)|^2 dx dy\right)^2}{\iint |E(x,y)|^4 dx dy}$$

where E(x, y) is the transverse electric field of the mode. The crosssection profile of TE mode field distribution for the coupling area of the resonator is shown in Fig. 2(b), in which about 1.363% of the power is confined in the slot between the waveguide and the microring.



Fig. 1. Schematic configuration of the Ge<sub>28</sub>Sb<sub>12</sub>Se<sub>60</sub> sensor based on an add-drop MRR: (a) 3D view, (b) Cross-section view of the coupling region, (c) Effective RI of the different order of TE/TM modes with the waveguide width at 1550 nm wavelength.



**Fig. 2.** Electric field distributions of (a) Fundamental TE mode in a 600 nm  $\times$  300 nm strip Ge<sub>28</sub>Sb<sub>12</sub>Se<sub>60</sub> waveguide with air upper cladding and SiO<sub>2</sub> substrate. (b) Cross-section field distribution of TE mode for the coupling area of the resonator. (c) Simulated bending loss variation as the radius of the microring resonator. (d) Simulated coupling length variation as the gap (the coupling coefficient is 0.01).

As mentioned before, the resonator is in a resonance state when the wavelength meets  $2\pi Rn_{eff} = m\lambda$ . A larger radius can increase the *Q* value of the sensor, thereby obtaining a lower detection limit. According to  $FSR = \lambda^2/(n_g \times 2\pi R)$ , a relatively small radius is needed to increase the sensing range. However, the bending loss will be increased sharply in a resonator with a very small radius. Fig. 2(c) illustrates the simulated results of the bending loss of a ring at different radii, where the bending loss is negligible when *R* is greater than 5 µm, and the radius of the Ge–Sb–Se MRR in our design is set as 24.6 µm. Fig. 2(d) also shows that when the power coupling coefficient is 0.01, coupling length will be increased by 50 times as the gap increases from 200 nm to 800 nm.

The transmission spectra of the designed MRR are also achieved based on 2.5D varFDTD solver (MODE Solutions). The transmission of the through and drop ports are shown in black and red lines in Fig. 3(a), respectively. The free spectrum range is calculated to be about 4.64 nm around the wavelength of 1550 nm from Fig. 3(b), which is in line with our design target. When the wavelength is far from the resonance wavelength as shown in Fig. 3(c), most of the energy is directly output from the through port of the micro ring, and a small part of the energy is coupled into the micro ring. When the wavelength satisfies the resonance condition as shown in Fig. 4(d), the energy is coupled from the bus waveguide, and resonance enhancement occurs in the ring and finally output from the drop port. In this case, the gap between the resonator and the waveguide is set as 200 nm to attain a near-critical coupling operation.

#### 3. Microring resonator fabrication and characterization

The device fabrication begins with depositing a 300 nm thick  $Ge_{28}Sb_{12}Se_{60}$  thin film by thermal evaporation on a 4-inch silicon wafer with a 2 µm thick silicon dioxide layer. A 300 nm thick negative electron beam resist (ma-N 2403) is then spin-coated onto the  $Ge_{28}Sb_{12}Se_{60}$  film at a spinning speed of 4000 rpm. After being patterned by electron beam lithography (Raith Eline Plus) and developed in ma-D525 solution, an inductively coupled plasma (ICP) reactive ion etcher (Oxford Instruments plasma 100) was then used to etch the ChG core layer with a gas mixture of CF<sub>4</sub> and CHF<sub>3</sub>, and the details of etching process could be found elsewhere [34]. When removing the residual photoresist, the device was first bombarded with the gas mixture of CHF<sub>3</sub>/AR/O<sub>2</sub> for 20 s, then placed in 1-Methyl-2-pyrrolidone solution and shaken at 70 °C for 3 min. After being rinsed in deionized water and purged with nitrogen, it was finally dried on a hot plate at 110 °C for 10 min to



Fig. 3. (a) Transmission spectra of through port (black) and drop port (red) of the microring resonator at the wavelength of 1500–1600 nm. (b) Partially enlarged image of micro ring transmission. The electric field distribution of the microring resonator at a position (c) far from the resonance wavelength and (d) a resonant wavelength position.

completely remove surface moisture. Fig. 4(a) shows the schematic fabrication of the Ge–Sb–Se micro ring sensor. Fig. 4(b)–(d) show the scanning electron micrographs (SEMs) of the fabricated device, and Fig. 4 (e) shows a fully etched TE type focused subwavelength grating coupler, which is fabricated to facilitate the coupling of light from the fiber into and out of the waveguide, with a period of 920 nm and a duty cycle of 0.715.

The schematic diagram of the setup for device performance characterization is shown in Fig. 5(a). A continuous-wave semiconductor laser with a wavelength-tunable range from 1500 nm to 1630 nm is used as the light source, and a manual fiber polarization controller is used to adjust the polarization state of the output beam. A single-mode fiber with tapered end and a fully etched TE mode grating is used to couple light into the sample. Then a CCD camera is used to monitor the position of the single-mode fiber and the sample, and alignment can be reached by a 3-axis translation stage. Finally, the output optical signal is connected to a high-precision and sensitivity power meter through a singlemode optical fiber and recorded by a computer.

Fig. 5(b) shows a typical transmission spectrum measured from the drop port of the microring resonator. The resonance is fitted using the Lorentzian function to extract a 3 dB linewidth of  $\delta \lambda = 0.04$  nm, as shown in Fig. 5(c). Using formula  $Q = \lambda_0 / \delta \lambda$ , the load Q factor of the resonator is calculated to be slightly higher than  $3.87 \times 10^4$ , and the

corresponding intrinsic *Q* factor can be estimated as  $7.74 \times 10^4$  according to  $Q_i = 2Q/(1 + \sqrt{T_0})$ , where  $T_0$  is the transmission dip at resonance. Finally, propagation loss of the microring is calculated to be 4.3 dB/cm.

The transmission spectrum of the device measured in different concentrations of NaCl solution is shown in Fig. 5(d), and the enlarged view of red dot lines area in (d) is illustrated in Fig. 5(e). It is obvious that the resonance peaks are red-shift as the solution concentration increases. For the NaCl aqueous solution at 20 °C, the refractive index increases by about 0.0018 for 1% increase in the concentration [35]. In our experiment, the concentration of NaCl solution increased from 0% to 5% with a 1% step, and its corresponding refractive index increased from 1.333 to 1.342. In order to minimize the influence of external temperature fluctuations on the sensing experiment, a thermoelectric cooler is used to maintain the temperature of the sample stage at 20 °C. Meanwhile, after each test, the device will be rinsed in deionized water and placed on a 110 °C hot plate to completely dry it.

Sensitivity and *LOD* are two important indexes for evaluating sensor performance. *LOD* of a sensor depends not only on its own performance, but also on system noise, light sources, detectors, and data processing methods. When the noise is negligible, the intrinsic limit of detection (*iLOD*) is employed to replace the system's limit [36]: *iLOD* =  $\lambda_{res}/(Q \cdot S)$ , where  $\lambda_{res}$  is the sensor's resonant wavelength, and *S* is the

(a)



Fig. 4. (a) Schematic fabrication of Ge–Sb–Se micro ring sensor. (b) Top view of the entire device. (c) Enlarged view of the coupling region. (d) SEMs of the waveguide cross-section. (e) Fully etched TE type focused grating coupler.

sensitivity. After linear fitting, as shown in Fig. 5(f), the sensitivity *S* is determined to be approximately 123.68 nm/RIU and the corresponding *iLOD* is about  $3.24 \times 10^{-4}$  RIU.

A comparison between various photonic integrated RI sensors employing microring configuration in recent years is presented in Table 1 [37–48], which provides a perspective on the features and benefits of the proposed device. The optical field of the SOI platformbased micro ring is highly confined in the waveguide core and cannot fully interact with external analytes, which will result in low sensitivity usually around 60 nm/RIU. Subwavelength grating microring and photonic crystal microring resonators are proposed to improve the sensitivity of the sensor further. By periodically etching sub hundred nanometer cavities or slots on the micro ring, the light field can more fully interact with the external analyte, such that the sensitivity can be increased to near 400 nm/RIU. However, this will not only greatly increase the design and fabrication complexity but also increase the transmission loss of the waveguide to tens of dB/cm, substantially reducing the *Q* factor of the device to the order of  $10^3$ . Several researchers have proposed to deposit a metal film on a dielectric to improve sensitivity by using the surface plasmon polariton effect at the



**Fig. 5.** (a) Schematic diagram of the experimental setup for transmission spectrum measurement. (b) The optical transmission spectrum of  $Ge_{28}Sb_{12}Se_{60}$  MRR of the TE mode at drop port. (c) Experimental data of spectrum (black points) and the Lorentz fitting (red solid line) of one resonant peak. (d) Measured transmission responses of MRR immersed in aqueous NaCl solution with different concentrations. (e) The enlarged view of red dot area in (d). (f) Resonant wavelength of microring sensor as a function of background RI. The sensitivity of the microring sensor is 123.68 nm/RIU by the linear fitting.

Table 1						
Comparison of	the proposed	microring sense	or and othe	er reported	microring	sensors

Year	Material	Туре	Radius (µm)	Q-factor	ER (dB)	Sensitivity (nm/RIU)	Exp or Sim
2014 [37]	SOI	Add-drop microring	5	$1.5 imes10^5$	16	65	simulation
2015 [38]	Si <sub>3</sub> N <sub>4</sub>	Differential microring	70	$1.9 imes10^5$	none	64.2	experiment
2015 [39]	Silicon + Silver	Hybrid plasmonic microring	5	$6.0  imes 10^2$	none	497	simulation
2016 [40]	SiON	Racetrack microring	100	$1.3 imes 10^5$	9	80	experiment
2017 [41]	SOI	Concentric dual-microring	5.5	none	30	180	simulation
2017 [42]	SOI	Photonic crystal microring	7.16	$1.2 imes 10^3$	10	248	experiment
2017 [43]	SOI	Subwavelength grating microring	10	$9.8  imes 10^3$	24.6	429.7	experiment
2018 [44]	SOI	Asymmetrical microring	1.7	$1.7  imes 10^3$	none	80	simulation
2018 [45]	Si <sub>3</sub> N <sub>4</sub>	All-pass microring	none	$5.0\times10^3\sim8.0\times10^3$	~8	72.5	experiment
2018 [46]	SOI	Phase-shifted Bragg grating microring	3.5	$2.0 imes 10^3$	20	297.13	simulation
2019 [47]	SOI	Subwavelength grating microring	10	$3.5 imes10^4$	none	672.8	simulation
2019 [48]	$Al_2O_3$	All-pass microring	200	$6.0  imes 10^5$	none	102.3	experiment
This work	$Ge_{28}Sb_{12}Se_{60}$	Add-drop microring	24.6	$7.7 imes10^4$	~40	123.68	experiment

interface of the two materials, but this solution encounters the same problem due to the inherent high absorption of metals. Other researchers have focused on materials with lower RI than silicon, such as  $Si_3N_4$ , SiON,  $Al_2O_3$ , and polymers. The decrease in RI means that the confinement of the optical field in the waveguide core materials is

weakened, such that the interaction between the optical field and the substance is strengthened to improve the sensitivity. But these micro rings need a much larger radius to avoid suffering from bending loss, which will be less compact for the structure of the device. Compared with the materials mentioned above, the moderate RI of  $Ge_{28}Sb_{12}Se_{60}$ 

ChG allows it to obtain an increase in sensitivity (compared with silicon) while maintaining a stronger confinement on the light field (compared with  $Al_2O_3$ , polymer, and  $Si_3N_4$ ) to avoid suffering from bending loss under a small bending radius, which will facilitate large-scale on-chip integration with other photonic devices.

#### 4. Conclusion

In this article, a compact, fabrication-friendly liquid sensor based on Ge<sub>28</sub>Sb<sub>12</sub>Se<sub>60</sub> MRR is proposed. First, the single-mode transmission conditions of the Ge<sub>28</sub>Sb<sub>12</sub>Se<sub>60</sub> waveguide are investigated at 1550 nm wavelength to determine the geometrical dimensions. Then, the TE mode profile of the waveguide and coupling region as well as the bending loss under different radii and coupling length under different gap widths are simulated. Finally, the sensor is fabricated based on electron beam lithography and tested under a CW tunable laser and a high-resolution power meter to obtain intrinsic O factors up to 7.74  $\times$ 10<sup>4</sup>. Device size could be reduced to the order of tens of microns without suffering from substantial bending loss due to the high RI contrast between the ChG core layer and SiO<sub>2</sub> substrate. Considering the relatively lower RI of Ge<sub>28</sub>Sb<sub>12</sub>Se<sub>60</sub> ChGs compared with silicon, a higher sensitivity is expected, and 123.68 nm/RIU is experimentally demonstrated, with an *iLOD* of  $3.24 \times 10^{-4}$ . Our research suggests that Ge<sub>28</sub>Sb<sub>12</sub>Se<sub>60</sub> chalcogenide materials have a great application prospect in sensing fields with high sensitivity.

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#### **Declaration of Competing Interest**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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